

intersecting D-brane model

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Ingredients

[quantum field theory](#)

[sigma-model](#),

[CFT](#), [perturbation theory](#)

[effective background QFT](#)

[gravity](#), [supergravity](#), [Yang-Mills theory](#), [quantum gravity](#)

Critical string models

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1. Idea

In [string phenomenology](#) for [type IIA string theory](#), quasi-realistic [models](#) (i.e. close to the [standard model of particle physics](#), or an [MSSM](#)) may be obtained by [KK-compactifications](#) from the [10-dimensions](#) down to 4d with configurations of [D6-branes](#) which fill all of 4d spacetime and [intersect](#) in a certain way in the 6-dimensional [fiber](#) space.

[U-duality](#)[open/closed string duality](#)[AdS/CFT correspondence](#)[holographic principle](#)[KLT relations](#)[11-dimensional supergravity](#), [M-theory](#)[Hořava-Witten theory](#)

Extended objects

[brane](#)

D-brane

[D0-brane](#), [D2-brane](#), [D4-brane](#)[D1-brane](#), [D3-brane](#), [D5-brane](#)[RR-field](#), [differential K-theory](#)

NS-brane

[string](#), [sigma-model](#)[spinning string](#), [superstring](#)[B2-field](#)[NS5-brane](#)[B6-field](#)

M-brane

[M2-brane](#)[C3-field](#)[ABJM theory](#), [BLG model](#)[M5-brane](#)[C6-field](#)[6d \(2,0\)-supersymmetric QFT](#)

Topological strings

[topological string](#), [TCFT](#)[A-model](#), [B-model](#)[topological M-theory](#)[Backgrounds](#)

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Fields and quanta

[fields](#) and [particles](#) in [particle physics](#)

and in the [standard model of particle physics](#):

[force field gauge bosons](#)

[photon](#) - [electromagnetic field](#)
([abelian Yang-Mills field](#))

[W, Z, B-boson](#) - [electroweak field](#)
([Yang-Mills field](#))

[gluon](#) - [strong nuclear force](#) ([Yang-Mills field](#))

[graviton](#) - [field of gravity](#)

[infraparticle](#)

[scalar bosons](#)

[Higgs boson](#), [inflaton](#) ([scalar field](#))

[matter field fermions](#) ([spinors](#),
[Dirac fields](#))

<u>flavors of fundamental fermions in the standard model of particle physics:</u>			
<u>generation of fermions</u>	1st generation	2nd generation	3d gene
<u>quarks</u> (<i>q</i>)			
up-type	<u>up quark</u> (<i>u</i>)	<u>charm quark</u> (<i>c</i>)	<u>top quark</u> (<i>t</i>)
down-type	<u>down quark</u> (<i>d</i>)	<u>strange quark</u> (<i>s</i>)	<u>bottom quark</u> (<i>b</i>)
<u>leptons</u>			
charged	<u>electron</u>	<u>muon</u>	<u>tauon</u>
neutral	<u>electron neutrino</u>	<u>muon neutrino</u>	<u>tau neutrino</u>
<u>bound states:</u>			

<u>mesons</u>	<p><u>light mesons:</u> <u>pion</u> (ud) <u>ρ-meson</u> (ud) <u>ω-meson</u> (ud) <u>f1-meson</u> <u>a1-meson</u></p>	<p><u>strange-mesons:</u> <u>ϕ-meson</u> ($s\bar{s}$), <u>kaon, K^*-meson</u> (us, ds) <u>eta-meson</u> ($uu + dd + ss$) <u>charmed heavy mesons:</u> <u>D-meson</u> (uc, dc, sc) <u>J/ψ-meson</u> ($c\bar{c}$)</p>	<p><u>bott heavy mesons</u> <u>B-mesons</u> (qb) <u>Υ-mesons</u> ($b\bar{b}$)</p>
<u>baryons</u>	<p><u>nucleons:</u> <u>proton</u> (uud) <u>neutron</u> (udd)</p>		

(also: antiparticles)

effective particles

Goldstone bosons

hadrons (bound states of the above quarks)

[meson](#)

[scalar meson](#)

[\$\sigma\$ -meson](#)

[pion](#)

[kaon](#)

[D-meson](#)

[B-meson](#)

[vector meson](#)

[\$\omega\$ -meson](#)

[\$\rho\$ -meson](#)

[f1-meson](#)

[a1-meson](#)

[b1-meson](#)

[h1-meson](#)

[K*-meson](#)

[tensor meson](#)

[quarkonium](#)

[charmonium](#)

[\$\Upsilon\$ -meson](#)

[exotic meson](#)

[XYZ meson](#)

[tetraquark](#)

[baryon](#)

[nucleon](#)

[proton](#)

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[Lambda baryon](#)

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[***solitons***](#)

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[instanton](#)

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[leptoquark](#)

[Z'-boson](#)

[*minimally extended supersymmetric standard model*](#)

[*superpartners*](#)

bosinos:

[gravitino](#) - [field of supergravity](#)
([Rarita-Schwinger field](#))

[gaugino](#) - [super Yang-Mills field](#)

[gluino](#)

[sfermions](#):

[squark](#)

[*dark matter candidates*](#)

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[axion](#)

Exotica

[preon](#)

[graviphoton](#)

[dilaton](#)

[monopole](#)

[dual graviton](#)

[giant graviton](#)

[*auxiliary fields*](#)

[ghost field](#)

[antifield](#)

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Quantum field theory

[functorial quantum field theory](#)

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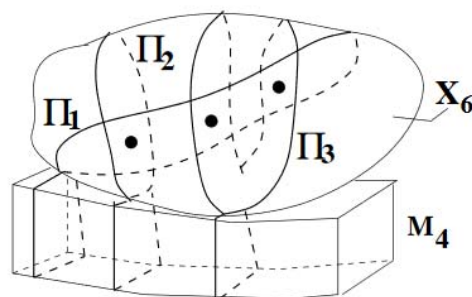


Figure 8: Compactification with intersecting D6-branes wrapped on 3-cycles.

graphics grabbed from [Uranga 12a, p. 12](#)

The [Chan-Paton gauge field](#) on the [D6-branes](#) yields the [gauge fields](#) in 4d, and the precise [intersection](#) pattern determines the effective [fundamental particle](#) content in 4d.

See also at [string phenomenology](#) the section [Models in type II with intersecting branes](#).

2. Bottom-up and Top-down approaches

One distinguishes [bottom-up and top-down model building](#) strategies:

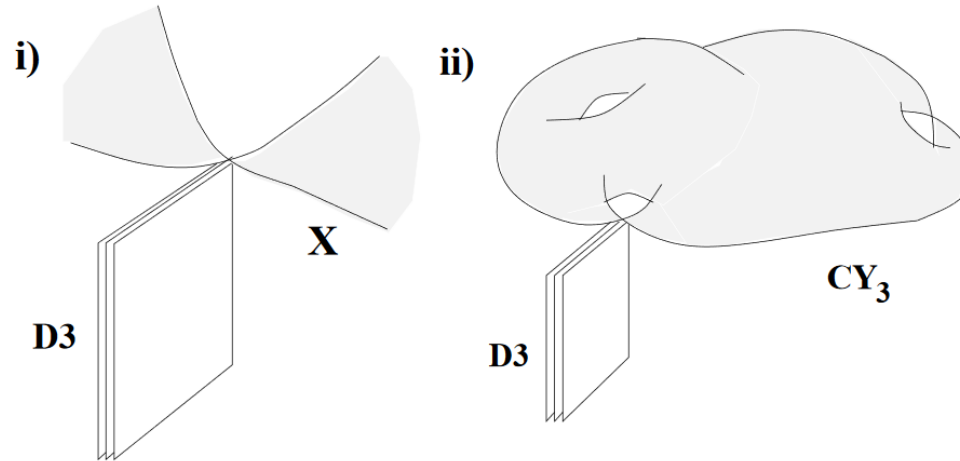


Figure 1: Pictorial representation of the bottom-up approach to the embedding of the standard model in string theory. In step **i)** the standard model is realized in the world-volume of D3-branes sitting at a singular non-compact space X (in the presence of D7-branes, not depicted in the figure). In step **ii)** this local configuration is embedded in a global context, like a compact Calabi-Yau threefold. Many interesting phenomenological issues depend essentially only on the local structure of X and are quite insensitive to the details of the compactification in step **ii)**. The global model may contain additional structures (like other branes or antibranes) not shown in the figure.

snippet grabbed from [Aldazabal-Ibáñez-Quevedo-Uranga 00](#)

3. Properties

Chiral fermions

One of the most striking special properties of the [standard model of particle physics](#) is that its content of [fermionic fundamental particles](#) is “chiral”, in that the left and right [Weyl spinor](#)-components of the would-be [Dirac spinor](#)-representation of the [quarks](#), [electrons](#) and [neutrinos](#) couple *differently* to the [gauge fields](#) (for review see e.g. [Ibanez-Uranga 12, section 1.1](#)).

The observation that such [chiral fermions](#) do indeed appear when [D6-branes intersect](#) at an angle on an $\mathbb{R}^{3,1}$ is due to [\(Berkooz-Douglas-Leigh 96\)](#), see also [\(AFIRU 00, section 4, BCLS 05, section 2.3\)](#). For review with an eye towards [RR-field tadpole cancellation on toroidal orientifolds](#) see also [Forste-Honecker-Schreyer 01, Honecker 01, Sec. 2, Honecker 02](#).

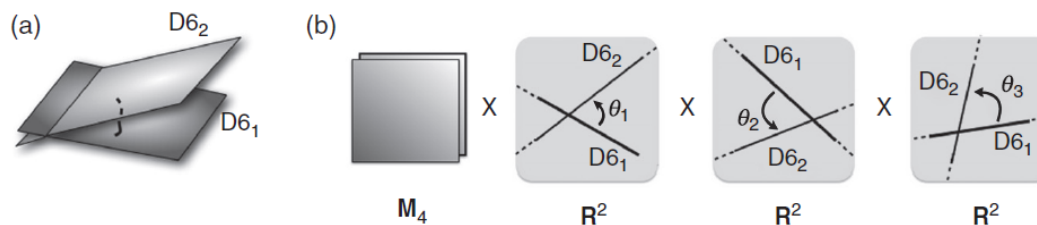


Figure 10.1 Two pictures of the configuration of two D6-branes intersecting over a 4d subspace of their volumes.

graphics grabbed from [Ibáñez-Uranga 12](#)

Vague review is in [\(Ibáñez-Uranga 12, section 10.2.1, Uranga 12a, section 2.3\)](#).

The lift of this situation to [M-theory on G2-manifolds](#) is discussed in [\(Berglund-Brandhuber 02, Bourjaily-Espahbodi 08\)](#).

Generations of fermions

While (presently) intersecting D-brane models don't explain why there are *precisely 3 generations of fundamental particles* in the [standard model of particle physics](#), they do have the property that generically any such model does feature several [generations of fundamental particles](#).

The reason is that in these models there is one copy of a set of fundamental particles at each intersection point of two 3-manifolds (the internal part of the D6-branes) in a compact 6-dimensional space, and generically these intersection numbers are greater than one and hence induce a finite number of [generations](#) [\(BCLS 05, section 2.3, Ibanez-Uranga 12, p. 307, Uranga 12a, p. 12\)](#):

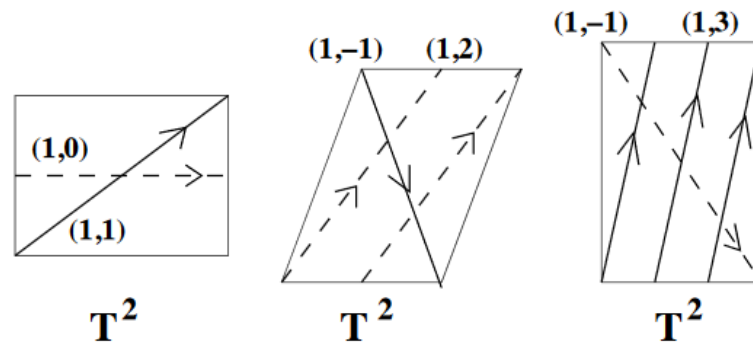


Figure 9: Examples of intersecting 3-cycles in T^6 .

graphics grabbed from [Uranga 12a, p. 13](#)

Yukawa couplings

In [intersecting D-brane models](#) [Yukawa couplings](#) are encoded by [worldsheet instantons](#) of open strings stretching between the [intersecting D-branes](#) (see [Marchesano 03, Section 7.5](#)). Mathematically this is encoded by [derived hom-spaces](#) in a [Fukaya category](#) (see [Marchesano 03, Section 7.5](#)).

A^∞ category	Fukaya Category	Intersecting D-branes
Objects	special Lagrangian submanifolds Π_α endowed with unitary systems \mathcal{E}_α	D-brane stacks with flat gauge bundles
Morphisms	$Hom(\mathcal{E}_\alpha, \mathcal{E}_\beta)$ generated by intersection points $p_{\alpha\beta} \in \Pi_\alpha \cap \Pi_\beta$	Chiral matter in bifundamentals localized at the intersections
Grading	Maslov index $\eta(p_{\alpha\beta})$	Spectral flow index (mod 2 = chirality)
Composition maps	Holomorphic discs m_1 m_2 $m_k, k \geq 3$	Worldsheet instanton corrections mass operator Yukawa coupling non-renormalizable couplings involving $k + 1$ chiral fields

Table 7.2: Fukaya-Yukawa dictionary.

table grabbed from [Marchesano 03](#)

Realistic [Yukawa couplings](#) and [fermion masses](#) in an [MSSM Pati-Salam GUT model](#) with 3 [generations of fermions](#) realized on [intersecting D6-branes KK-compactified](#) on a [toroidal orbifold](#) $T^6 // (\mathbb{Z}_2 \times \mathbb{Z}_2)$ are claimed in in [Mayes 19](#), [Gemill-Howington-Mayes 19](#), based on [Chen-Li-Mayes-Nanopoulos 07a](#), [Chen-Li-Mayes-Nanopoulos 07b](#).

Higgs mechanism

The [Higgs mechanism](#) naturally arises in [intersecting D-brane models](#): The [Higgs field](#) appears here as the [scalar field](#) that witnesses in [perturbation theory](#) the process of *brane recombination* at the intersection locus of the D-branes ([Cremades-Ibanez-Marchesano 02, section 7](#)):

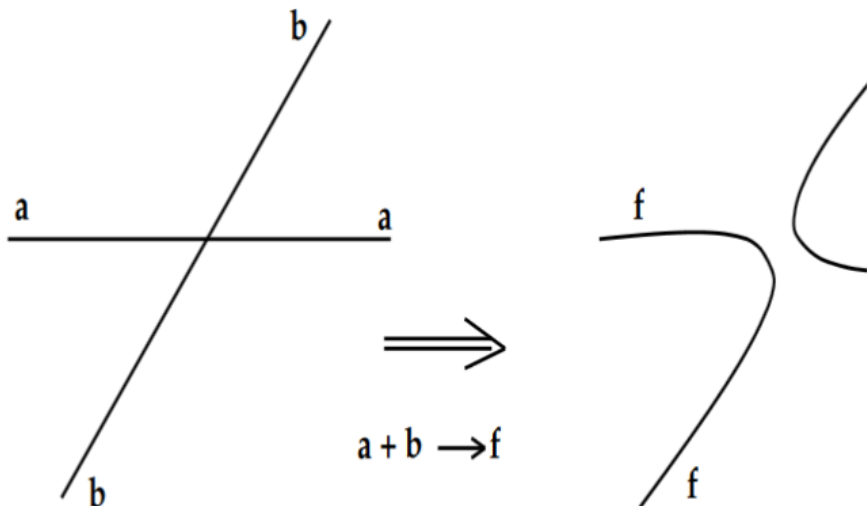


Figure 11: Recombination of two intersecting branes a, b in a compact space. Initially the gauge group is $U(1)_a \times U(1)_b$. At the intersection a tachyon scalar triggers the recombination into a single brane f . The final gauge symmetry is $U(1)_f$, corresponding to the Higgsing $U(1)_a \times U(1)_b \rightarrow U(1)_f$ which is induced by the tachyon.

Further developments in [Ibanez-Marchesano-Rabadan 01](#), [Hebecker-Knochel-Weigand 13](#) and specifically via [string field theory](#) in [HINSS 18](#).

For review see [Ibanez-Uranga 12, fig 10.2](#):

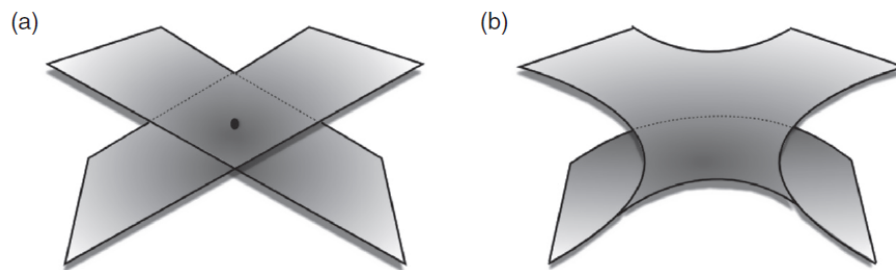


Figure 10.2 Recombination of two intersecting D6-branes into a single smooth one, corresponding to a vev for a scalar at the intersection.

RR-Tadpole cancellation and Orientifolding

Consistent intersecting D-brane models have to be in [type I string theory](#), or generally in [type II string theory](#) with [orientifold](#) backgrounds, to achieve [RR-field tadpole cancellation](#).

This is a key consistency condition in intersecting D-brane model building (e.g. [BCLS 05, section 2.4](#), [Ibanez-Uranga 12, section 4.4](#))

RR-field tadpole cancellation conditions for D-branes wrapped on toroidal orientifolds in terms of their [D-brane charge](#) $V \in K_G(*) = R_G$ in [equivariant K-theory](#) \equiv [representation ring](#) (here $\mathbf{n}_{\text{reg}} = k[G/1]$ denotes the [regular representation](#) of [dimension](#) $n = \text{ord}(G)$)

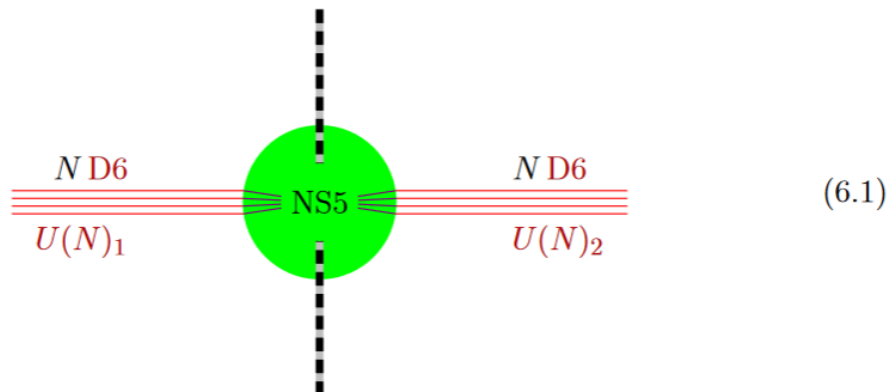
single D-brane species on toroidal orientifold	local/twisted tadpole cancellation condition	global/untwisted tadpole cancellation condition	comment
D5-branes transv. to $\mathbb{T}^{4\text{H}} // \mathbb{Z}_2$	$V = N \cdot \mathbf{2}_{\text{reg}}$ (Buchel-Shiu-Tye 99 (19))	$V = 16 \cdot \mathbf{2}_{\text{reg}}$ (Buchel-Shiu-Tye 99 (18))	following Gimon-Polchinski 96 , Gimon-Johnson 96
D5-branes transv. to $\mathbb{T}^{4\text{H}} // \mathbb{Z}_4$	$V = N \cdot \mathbf{4}_{\text{reg}}$ (Buchel-Shiu-Tye 99 (19))	$V = 8 \cdot \mathbf{4}_{\text{reg}}$ (Buchel-Shiu-Tye 99 (18))	following Gimon-Polchinski 96 , Gimon-Johnson 96
D4-branes transv. to $\mathbb{T}^{1\text{triv}+4\text{H}} // \mathbb{Z}_k$	$V = N \cdot \mathbf{k}_{\text{reg}}$ (AFIRU 00a, 4.2.1)		
D4-branes transv. to $\mathbb{T}^{1\text{triv}+4\text{H}} // \mathbb{Z}_3$	$V = N \cdot \mathbf{3}_{\text{reg}}$ (AFIRU 00b, (7.2))	$V = 4 \cdot \mathbf{3}_{\text{reg}} + 4 \cdot \mathbf{1}_{\text{triv}}$ (Kataoka-Shimojo 01, (14)-(17))	

single <u>D-brane</u> species on <u>toroidal orientifold</u>	<u>local/twisted tadpole cancellation condition</u>	<u>global/untwisted tadpole cancellation condition</u>	comment
<u>D8-branes</u> on $\mathbb{T}^{1_{\text{triv}}+4_{\mathbb{H}}} // \mathbb{Z}_3$	$V = N \cdot \mathbf{3}_{\text{reg}}$	$V = 4 \cdot \mathbf{3}_{\text{reg}} + 4 \cdot \mathbf{1}_{\text{triv}}$ (Honecker 01, 4 , Honecker 02a, (25) \wedge (28) \Leftrightarrow (29) , Honecker 02b, (3.19)-(3.27))	equivalent to D4 case by <u>T-duality</u> : Honecker 01, p. 2 , Honecker 02a, 6 , Honecker 02b, p. 15 review in: Marchesano 03, Sec. 4
<u>D6-branes</u> on $\mathbb{T}^6 // \mathbb{Z}_4$		$V = 8 \cdot \mathbf{4}_{\text{reg}}$ (Ishihara-Kataoka-Sato 99, (4.16))	
<u>D3-branes</u> on $\mathbb{T}^{4_{\mathbb{H}}} // \mathbb{Z}_k$	$V = N \cdot \mathbf{k}_{\text{reg}}$ (Feng-He-Karch-Uranga 01, (25))		
<u>D7-branes</u> on $\mathbb{T}^{4_{\mathbb{H}}} // \mathbb{Z}_k$	$V = N \cdot \mathbf{k}_{\text{reg}}$ (Feng-He-Karch-Uranga 01, (5), (6))		

Intersections of D6s with D8/O8s

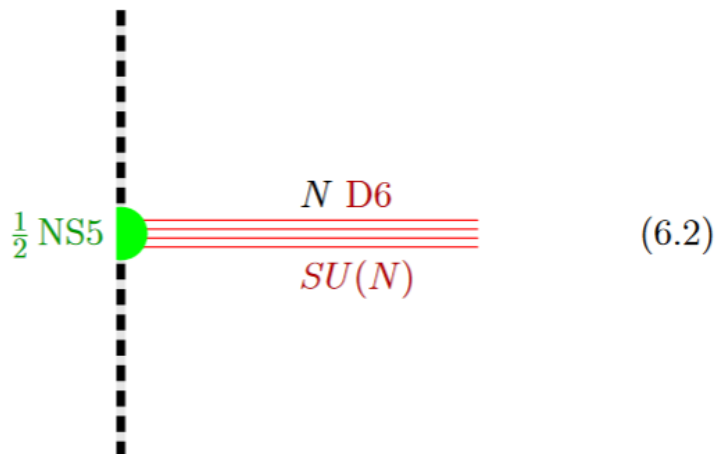
A black D6-brane may end on a black NS5-brane, and in fact each brane NS5-brane at a non-trivial ADE singularity has to be the junction of two black D6-branes. For details see at [D6-branes ending on NS5-branes](#).

An NS5 half-brane terminus of N D6 branes results from $O8$ orientifold projection of the following picture:



from [GKSTY 02](#)

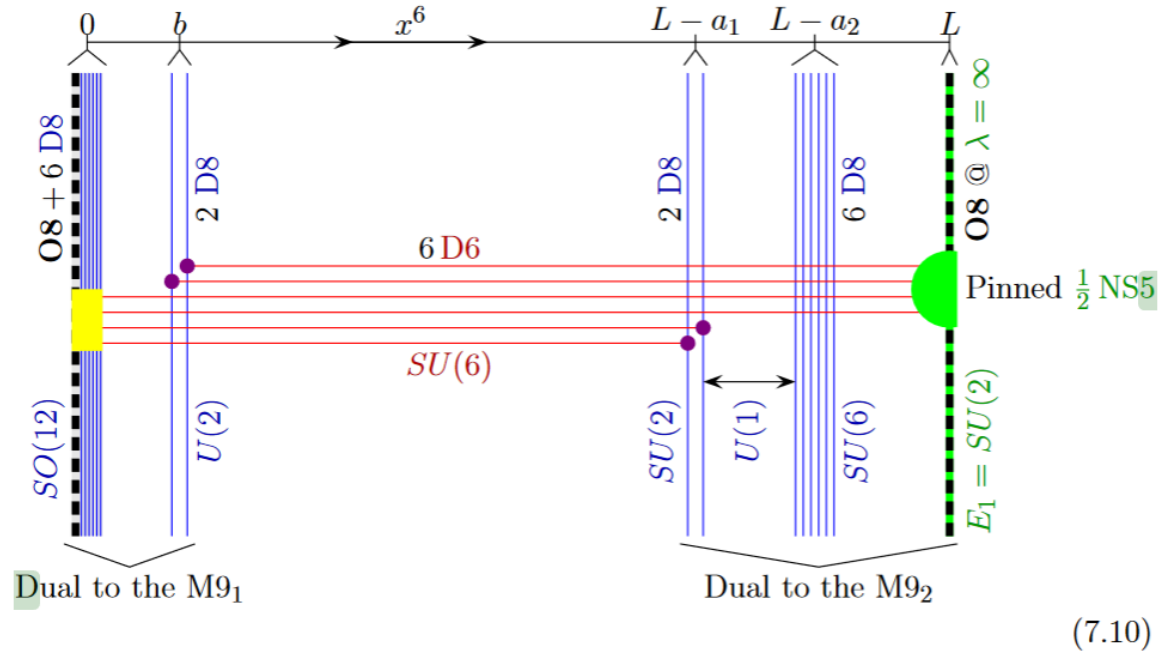
If in addition the [black NS5-brane](#) sits at an [O8-plane](#), hence at the [orientifold fixed point](#)-locus (see [above](#)), then in the ordinary $\mathbb{Z}/2$ -[quotient](#) it appears as a “[half-brane](#)” with only one copy of [D6-branes](#) ending on it:



from [GKSTY 02](#)

(In [Hanany-Zaffaroni 99](#) this is interpreted in terms of the '[t Hooft-Polyakov monopole](#).)

The lift to M-theory of this situation is an M5-brane intersecting an M9-brane:



from [GKSTY 02](#)

Alternatively, the O8-plane may intersect the black D6-branes away from the black NS5-brane:

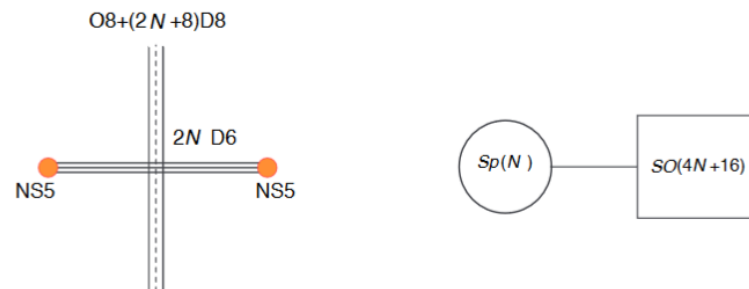


Figure 1. Left: Type IIA brane realization of the (D_{N+4}, D_{N+4}) minimal conformal matter in the tensor branch. Right: The quiver diagram of the 6d theory.

	0	1	2	3	4	5	6	7	8	9
D6-brane	×	×	×	×	×	×	×			
NS5-brane	×	×	×	×	×	×				
D8-brane	×	×	×	×	×	×		×	×	×
O8-plane	×	×	×	×	×	×		×	×	×

Table 1. The brane configuration that yields a 6d theory on a tensor branch of a 6d SCFT in Type IIA string theory.

from [HKLY 15](#)

In general, some of the NS5 sit away from the [O8-plane](#), while some sit on top of it:

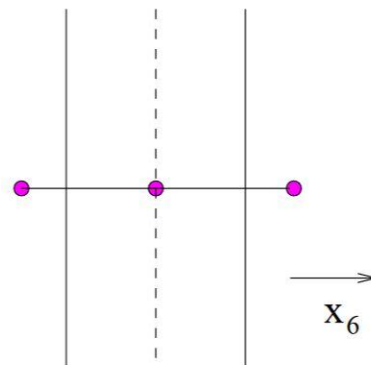


Figure 3: SU gauge group with a (anti)symmetric matter coming from D6-branes stretched between the NS-brane stuck to the orientifold $O8$ ($O8'$) plane and another NS-brane.

from [Hanany-Zaffaroni 98](#)

Relation to M-theory on G_2 -manifolds

Lift to [M-theory on \$G_2\$ -manifolds](#) (e.g. [\$G_2\$ -MSSM](#)): see references [below](#)

Cosmology and Holography

Since the [near horizon geometry](#) of [BPS black branes](#) is conformal to the [Cartesian product](#) of [anti de Sitter spaces](#) with the unit n -sphere around the brane, the [cosmology](#) of intersecting D-brane models realizes the [observable universe](#) on the [asymptotic boundary](#) of an [approximately anti de Sitter spacetime](#) (see for instance [Kaloper 04](#), [Flachi-Minamitsuji 09](#)). The basic structure is hence that of [Randall-Sundrum models](#), but details differ, such as notably in [warped throat](#) geometries, see [Uranga 05, section 18](#).

These warped throat models go back to [Klebanov-Strassler 00](#) which discusses aspects of [confinement](#) in [Yang-Mills theory](#) on coincident ordinary and [fractional D3-branes](#) at the [singularity](#) of a warped [conifold](#). See also [Klebanov-Witten 98](#).

flux. For $K \gg g_s M$, where g_s is the IIB string coupling, fluxes backreact on the geometry and create a **warped throat**, whose geometry is **approximately $AdS_5 \times T^{1,1}$** (with the cosmological constant of AdS_5 and the 5-form over $T^{1,1}$ slowly varying along the radial direction). The roles of the IR and UV branes in RS1 are played by the S^3 of the deformed conifold and the cutoff/compactification, respectively. The relative warp factors between both endpoints is $e^{-\frac{2\pi K}{3g_s M}}$ [73, 72, 43]. See figure 29

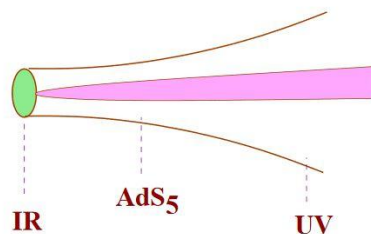


Figure 29: Schematic representation of the KS throat.

The holographic dual description is given by the **(almost conformal)** $\mathcal{N} = 1$ supersymmetric 4d gauge theory on $N = KM$ D3-branes at the singular conifold in the presence of M fractional branes. The gauge theory [74] has a gauge group $SU(N) \times$

snippet grabbed from [Uranga 05, section 18](#)

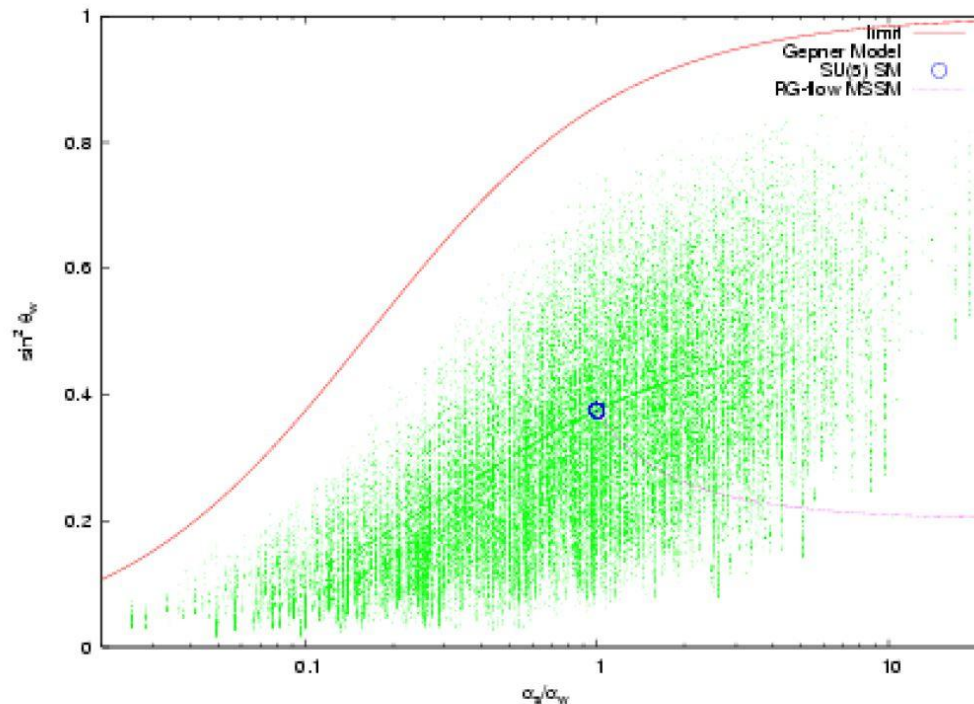
here: “RS”=[Randall-Sundrum model](#); “KS”=[Klebanov-Strassler 00](#)

In particular this means that [AdS-CFT duality](#) applies in *some approximation* to intersecting D-brane models (e.g. [Soda 10](#), [GHMO 16](#)), thus allowing to compute, to some approximation, [non-perturbative effects](#) in the [Yang-Mills theory](#) on the intersecting branes in terms of [gravity](#) on the ambient warped throat \sim [AdS](#) ([Klebanov-Strassler 00, section 6](#))

Such approximate versions of [AdS-CFT](#) for gauge theories realized on intersecting D-branes are used for instance to estimate [non-perturbative effects](#) in [QCD](#), such as the [shear viscosity](#) of the [quark-gluon plasma](#) (see the references [there](#)). See at [AdS-QCD correspondence](#) for more on this.

Computer scan of Gepner-model compactifications

Discussion of [string phenomenology](#) of [intersecting D-brane models](#) [KK-compactified](#) with non-geometric [fibers](#) such that the would-be string [sigma-models](#) with these [target spaces](#) are in fact [Gepner models](#) (in the sense of [Spectral Standard Model and String Compactifications](#)) is in ([Dijkstra-Huiszoon-Schellekens 04a](#), [Dijkstra-Huiszoon-Schellekens 04b](#)):



[65]. Approximately 10% of the models is on that line. Not surprisingly, these models are characterized by the fact that the reflection coefficients for respectively branes a and d and branes b and c are the same, which is precisely the RCFT equivalent of the condition that [76] used for unification. So although a portion of the models we found has a relation between the coupling constants compatible with $SU(5)$ unification, this is certainly not a

A plot of [standard model-like coupling constants](#) in a computer scan of [Gepner model-KK-compactification](#) of [intersecting D-brane models](#) according to [Dijkstra-Huiszoon-Schellekens 04b](#).

The blue dot indicates the couplings in $SU(5)$ -[GUT](#) theory. The faint lines are NOT drawn by hand, but reflect increased density of Gepner models as seen by the computer scan.

4. Examples

General

brane intersections/bound states/wrapped branes/polarized branes

- D-branes and anti D-branes form bound states by tachyon condensation, thought to imply the classification of D-brane charge by K-theory.
- intersecting D-branes/fuzzy funnels:
 - Dp-D(p+2)-brane bound state and Yang-Mills monopoles
 - D0-D2 brane bound state
 - D1-D3 brane bound state
 - D2-D4 brane bound state
 - D3-D5 brane bound state
 - D4-D6 brane bound state
 - D6-D8-brane bound state
 - Dp-D(p+4)-brane bound state and Yang-Mills instantons
 - D0-D4 brane bound state
 - D1-D5 brane bound state
 - D2-D6 brane bound state
 - D3-D7 brane bound state
 - D4-D8 brane bound state
 - D5-D9 brane bound state
small instanton
- Dp-D(p+6) brane bound state
 - D0-D6 brane bound state
 - D1-D7 brane bound state
 - D2-D8 brane bound state
- NS5-D4-D2 bound states

S-duality bound states:

- (p,q)-string
- (p,q)5-branes
- string network
- 5-brane web

intersecting M-branes:

- M5-M09 brane bound state
- M2-M5 brane bound state
 - M-string, E-string
- membrane triple junction

Intersecting D6-brane models

(...)

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(...)

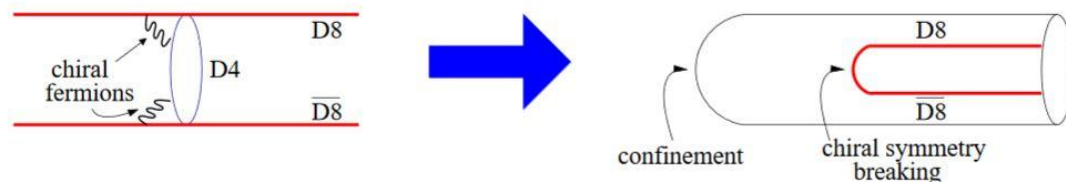
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D4-D8 brane bound states:

give the Witten-Sakai-Sugimoto model for holographic QCD



graphics grabbed from [Erlich 09, section 1.1](#)

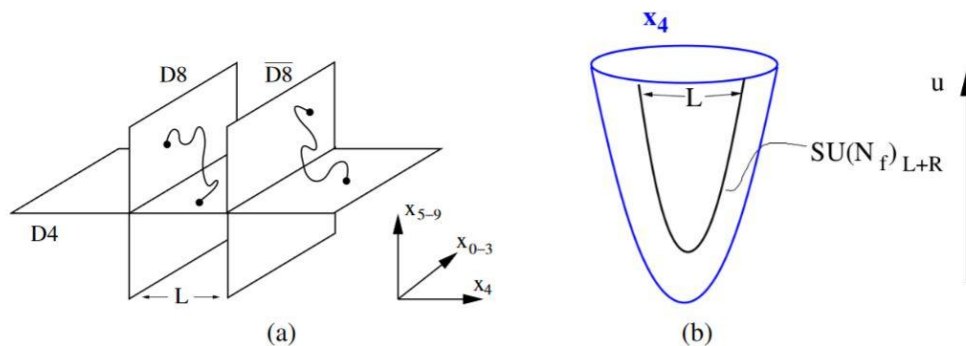


Figure 2. (a) D4-D8-D8 brane configuration in the 10 dimensions x_0, \dots, x_9 , with x_4 understood as periodic; (b) near-horizon geometry with u being a radial coordinate in the transverse space $x_{5\dots 9}$.

graphics grabbed from [Rebhan 14](#)

(...)

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